

# Secrets of Transmission Lines

## Part 6: The Smith Chart.

*In the last chapter, we looked at the behavior of transmission lines under steady-state conditions, and paid particular attention to the variation of impedance along the line for various termination conditions. In the experiment, we noted the repetition of the termination impedance every half-wave as well as the cycling of the impedance between  $Z_0 \cdot VSWR$  and  $Z_0/VSWR$ . In this chapter, we are going to look at what happens between the pure resistance points.*

In the previous work, we showed the expressions for the voltage and current along the line. From these expressions previously given, you can obtain the expression:

$$Z_x = E_x/i_x$$

eqn (6-1)

This simply says that at a point  $x$  along the line, the impedance is given by the voltage at that point divided by the current at that point. Bear in mind that both voltage and current can have forward and reflected components and may have real and imaginary parts.

The previously stated expressions for  $E_x$  and  $i_x$  can be substituted into this expression to develop the equation for  $Z_x$ . The actual manipulation is too long-winded for this treatment; however, for those interested, my version of it may be found in *Exploring Antennas and Transmission Lines by Personal Computer*, published by Van Nostrand Reinhold, New York. This book is now out of print, but still is in the possession of a number of hams and libraries. Other more current texts also carry the discussion.

As another shorthand notation to

simplify the printing of the equations we define

$$\beta = 2\pi/\lambda$$

(6-2)

where  
 $\lambda$  = wavelength in the medium

It is also common to refer to the position  $j$  along the line as " $\ell$ " rather than " $x$ " as we have been doing; however, for the purposes of the computer program to follow, we will retain the " $x$ ".

Using these conventions, we may write:

$$Z_x = Z_0 \cdot \left\{ \frac{[Z_1 \cdot \cos(\beta x)] + [j \cdot Z_0 \cdot \sin(\beta x)]}{[Z_0 \cdot \cos(\beta x)] + [j \cdot Z_1 \cdot \sin(\beta x)]} \right\}$$

(6-3)

where  
 $Z_1$  is the terminating impedance  
 $Z_x$  is as defined in (6-1)  
 $\beta$  is as defined in (6-2)

This is the complete expression for the impedance at any point " $x$ " on the transmission line. Note that it has real and imaginary (reactive; remember

chapter 2!) parts and that the terminating element  $Z_i$  can have real and imaginary parts as well.

For a particularly interesting case, let us set  $Z_i = 0$  (that is, the end of the line is perfectly short-circuited). For the shorted line  $Z_i = 0$ , the upper left and lower right terms are zero. Therefore

$$Z_x = Z_0 \cdot \left\{ \frac{[0 + [j \cdot Z_0 \cdot \sin(\beta x)]]}{[Z_0 \cdot \cos(\beta x)] + 0} \right\}$$

(6-4)

$$Z_x = j \cdot Z_0 \cdot \tan(\beta x)$$

(6-5)

Equation (6-5) follows from the fact that  $\sin(a)/\cos(a) = \tan(a)$ .

The tangent function is such that if  $(b \cdot x) = 45$  degrees, then  $Z_x = j \cdot Z_0$ . In other words, a shorted section of line an eighth of a wave long behaves like an inductor with a reactance equal to  $Z_0$  ohms. At the half-wave point, the value of the tangent goes to infinity and the line section or stub looks like an open circuit, as we noted with the experiment in the previous chapter. As a matter of fact, the tangent changes algebraic sign just beyond a half wave and the stub looks like a parallel



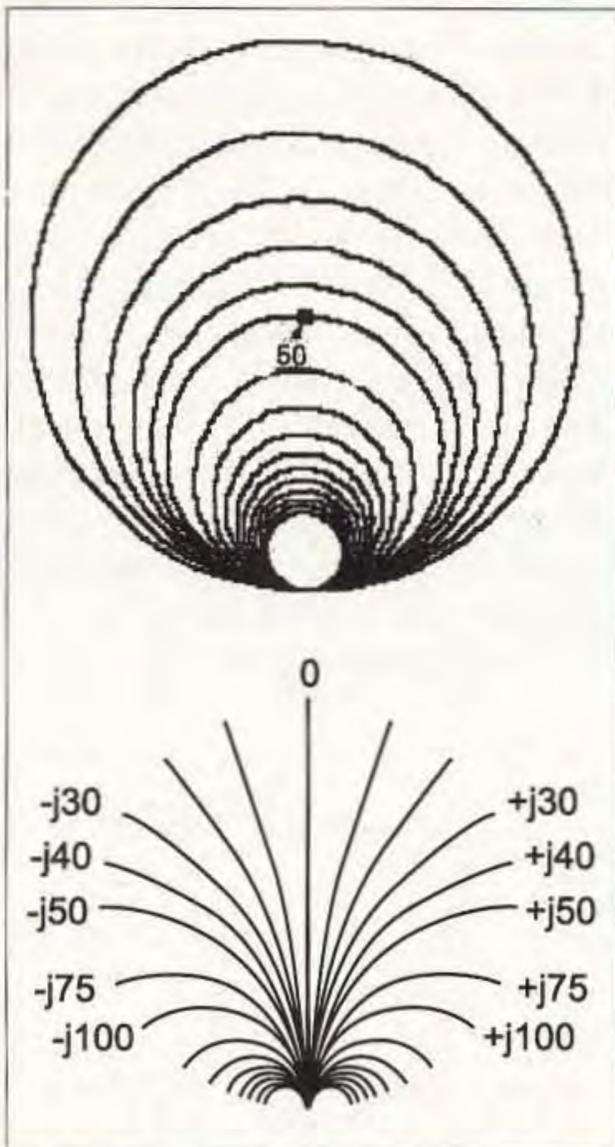


Fig. 2. The Smith Chart makeup. Top: Resistance curves are shown in 10-ohm steps for  $R < 50$  and 25-ohm steps for  $R > 50$ . Bottom: Reactance curves.

impedances the line goes through. A quarter wave on the transmission line represents a half turn on the chart; the full circle represents a half wave, and the impedance repeats itself just as we saw in the last chapter.

Just think about this for the moment. By simply drawing the VSWR circle, we solved the transmission line equations for that load or termination for all possible line lengths.

Fig. 2 shows some of the resistance circles and some of the reactive circles on the chart. You will note that the zero reactance curve is the centerline and the reactance has a non-zero value everywhere else. The resistance circles and the reactance circles are said to be orthogonal, meaning that they always cross at right angles. Also note that all of the circles pass through the  $R/Z_0 = \text{infinity}$  point.

Fig. 3 shows how the circles are generated. Looking at Fig. 3(a), we can see that if we terminate the line in a short circuit at  $R/Z_0 = \text{zero}$ , then a quarter of a wavelength down the

line we will have  $R/Z_0 = \text{infinity}$ , and another quarter wave takes us back to the zero point. The  $R = 0$  circle is the outer periphery of the chart. All possible impedances with real and reactive parts ranging from zero to

infinity can be plotted on the Smith chart.

Also given on Fig. 3 are the formulas for generating a Smith Chart in admittance terms. It is often convenient to work in admittance terms, since it is



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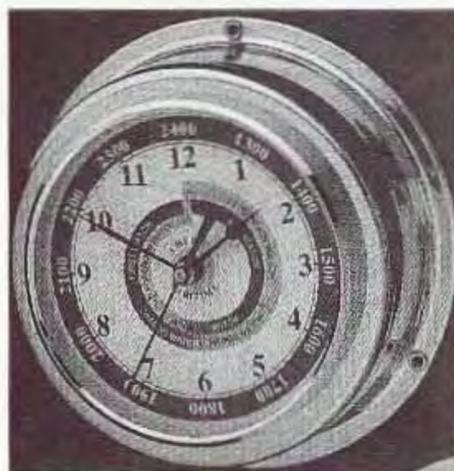
usually easier to place corrective or impedance matching elements in shunt across a coaxial cable rather than in series. From a practical standpoint, you can cut the cable and install a "tee" and hang a shunt element there more easily than you can insert a series element. The

reciprocal of 50 ohms is 0.020 mhos or siemens.

### An example of impedance matching with the Smith Chart

Next, let us consider an example of impedance matching using the Smith Chart. We will work in admittance because I intend to do the matching with a short-circuited

stub of line of the same  $Y_0$  ( $1/Z_0$ ) as the line. The example is shown in Fig. 4. We start with a chart having a 20 millimho center and plot the load on it, which is given as  $10 + j2$  mmhos. Next, we draw the VSWR circle. Only part of it is shown, for clarity. We rotate the arc until it meets the 20 mmho circle. The rotation is clockwise toward the generator and counterclockwise away from the generator. Next, we lay a straightedge from the center to the circular scale on the outside of the chart. The original load point reads .02 wavelengths and the point where



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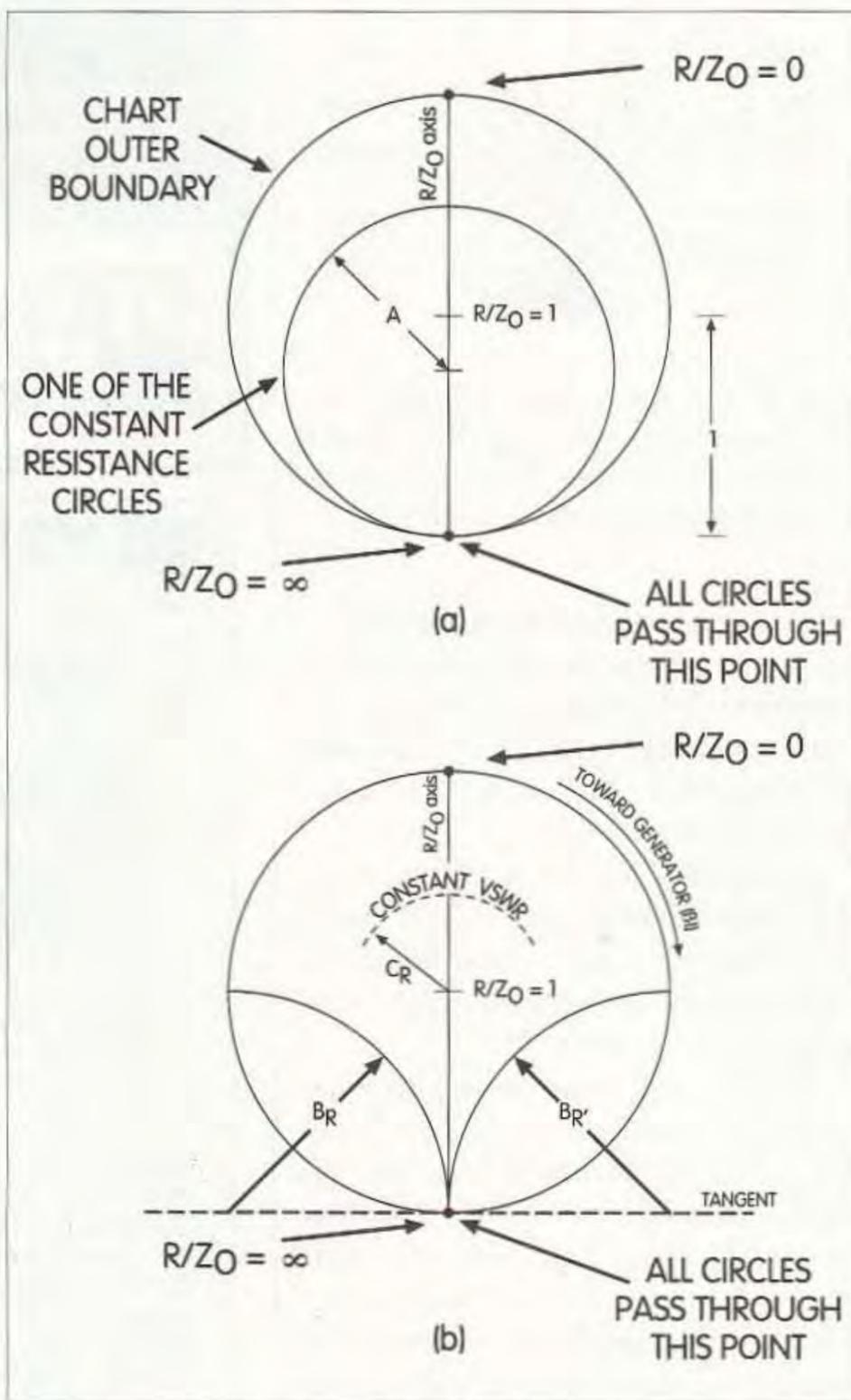


Fig. 3. Further makeup of the Smith Chart.

$$A = \frac{1}{1 + \frac{R}{Z_0}} \text{ or } A = \frac{1}{1 + \frac{G}{Y_0}} \quad C_R = 2A - 1 = \frac{2}{1 + \frac{R}{Z_0}} - 1$$

$$B_R = \frac{Z_0}{X} \text{ or } B_R = \frac{Y_0}{B}$$



# Secrets of Transmission Lines

*continued from page 35*

around the twinlead. This makes a capacitor. See whether you can find a location along the twinlead where you can make the line impedance match or flatten out on the generator side. Try the same trick with a 600-ohm termination.

## Conclusion

Next time, we will conclude the series and present some computer programs suited to transmission line work and impedance matching, as well as have a general discussion on which circuits are appropriate for which impedances.